CIEM5000: Structural Engineering Base Matrix Method – Final Details

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The Matrix Method

Main steps:

- Extract element matrices
- Impose nodal equilibrium
- Impose boundary conditions
- Solve for unknown displacements
- Postprocess results

This week:

- Element loads
- Non-zero Dirichlet boundary conditions in two different ways
- Postprocessing: support reactions and element fields
- Matrix method versus FEM parallels and differences
- Example: Support reactions of extension bar with distributed load
- Example: A fully-resolved example by hand
- Workshop: Wrap up the code and solve a frame structure

Element loads

The matrix method is a discrete approach

- Nodal loads treated easily
- What if we have loads applied inside elements?

A number of approaches to handle this:

- Further discretization (seldom helps)
- ODE approach
- Work-based approach



We can follow the same steps as before: General solution for the ODE:

$$EA\frac{\mathrm{d}^2 u}{\mathrm{d}x^2} = -q \quad \Rightarrow \quad u(x) = -\frac{qx^2}{2EA} + C_1 x + C_2$$



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$$u(0) = u_1 \quad u(\ell) = u_2 \quad \Rightarrow \quad C_1 = \frac{q\ell}{2EA} + \frac{u_2 - u_1}{\ell} \quad C_2 = u_1$$
$$u(x) = \frac{q}{2EA} \left(\ell x - x^2\right) + u_1 \left(1 - \frac{x}{\ell}\right) + \frac{u_2 x}{\ell}$$

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Relate forces at the edges with internal stresses:

$$N(x) = \frac{EA}{\ell} (u_2 - u_1) + \frac{q\ell - 2qx}{2} \qquad F_1 = -N_1 \quad F_2 = N_2$$

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Relate forces and displacements, but now an extra term appears:

$$\frac{EA}{\ell} \begin{bmatrix} 1 & -1 \\ -1 & 1 \end{bmatrix} \begin{bmatrix} u_1 \\ u_2 \end{bmatrix} - \begin{bmatrix} \frac{q\ell}{2} \\ \frac{q\ell}{2} \end{bmatrix} = \begin{bmatrix} F_1 \\ F_2 \end{bmatrix}$$

Dealing with equivalent loads

The new term is an equivalent nodal load:

- Element loads \Rightarrow nodal loads
- Force equilibrium at the nodes therefore changes a bit:

$$-\sum_{e} \mathbf{f}^{e} + \mathbf{f}_{\text{nodal}} = \mathbf{0}$$

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Remember, this is a force in the global coordinate system!

$$\mathbf{f}_{\mathrm{eq}} = \mathbf{T}^{\mathrm{T}} \overline{\mathbf{f}}_{\mathrm{eq}}$$



$\stackrel{H}{\longrightarrow} \underset{1}{\overset{F_1^{(1)}}{\longleftarrow}} (1) \stackrel{F_2^{(1)}}{\overset{F_2^{(2)}}{\longleftarrow}} \stackrel{F_1^{(2)}}{\overset{F_2^{(2)}}{\longleftarrow}} (2) \stackrel{F_2^{(2)}}{\overset{F}{\longleftarrow}} \stackrel{F}{\underset{3}{\longrightarrow}}$

Euler-Bernoulli bending, point load at midspan:

• Displacement field for arbitrary DOFs (ODE with q = 0):

$$w(x) = \underbrace{\left(\frac{2x^3}{\ell^3} - \frac{3x^2}{\ell^2} + 1\right)}_{s_1} w_1 + \underbrace{\left(-\frac{x^3}{\ell^2} + \frac{2x^2}{\ell} - x\right)}_{s_2} \varphi_1 + \underbrace{\left(-\frac{2x^3}{\ell^3} + \frac{3x^2}{\ell^2}\right)}_{s_3} w_2 + \underbrace{\left(-\frac{x^3}{\ell^2} + \frac{x^2}{\ell}\right)}_{s_4} \varphi_2$$

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EL.

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 EI, ℓ



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• Work performed by edge forces:

$$W_F = F_1^{\text{eq}} w_1 + T_1^{\text{eq}} \varphi_1 + F_2^{\text{eq}} w_2 + T_2^{\text{eq}} \varphi_2$$

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 $W_q = Pw(\ell/2) = (Ps_1(\ell/2)) w_1 + (Ps_2(\ell/2)) \varphi_1 + (Ps_3(\ell/2)) w_2 + (Ps_4(\ell/2)) \varphi_2$

 $T_{q}^{q} \underbrace{ \begin{array}{c} \varphi_{I} \\ \varphi_{I} \end{array}}_{EI, \ell} \underbrace{ \begin{array}{c} \varphi_{2} \\ \varphi_{2} \end{array}}_{EI, \ell} \underbrace{ \begin{array}{c} \varphi_{2} \\ \varphi_{2} \end{array}}_{T_{2}^{q}} \underbrace{ \begin{array}{c} \varphi_{I} \\ \varphi_{I} \end{array}}_{T_{2}^{q}} \underbrace{ \begin{array}{c} \varphi_{I} \end{array}}_{T_{2}^{q}} \underbrace{ \end{array}}_{T_{2}^{q}} \underbrace{ \begin{array}{c} \varphi_{I} \end{array}}_{T_$

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 $\overset{T^{rq}}{\overbrace{}} \overbrace{\overset{\varphi_{I}}{\overbrace{}}}^{\varphi_{I}} \overset{\varphi_{2}}{\underset{EI, \ell}{\overbrace{}}} \overbrace{\overset{\varphi_{2}}{\overbrace{}}}^{\varphi_{2}} \overbrace{\overset{\varphi_{2}}{\overbrace{}}}^{T^{rg}}$

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G

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After solving, we can easily recover support reactions:

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 \blacksquare Note: \mathbf{f}_{c} includes both nodal loads, equivalent loads and support reactions.

Dirichlet BCs - Size-preserving approach

The approach from before can be annoying to code:

- Reordering the system costs computation time
- Gains when inverting the stiffness matrix are very limited $(N_c \ll N_f)$

Alternatively, we can modify the relevant equations and solve the full system:

Support reactions recovered later from the unconstrained system

K_{11}	K_{12}	K_{13}	$\begin{bmatrix} u_1 \end{bmatrix}$		f_1
K_{21}	K_{22}	K_{23}	u_2	=	f_2
K_{31}	K_{32}	K_{33}	u_3		f_3

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$$\begin{bmatrix} K_{11} & 0 & K_{13} \\ 0 & 1 & 0 \\ K_{31} & 0 & K_{33} \end{bmatrix} \begin{bmatrix} u_1 \\ u_2 \\ u_3 \end{bmatrix} = \begin{bmatrix} f_1 - K_{12}\Delta_2 \\ \Delta_2 \\ f_3 - K_{32}\Delta_2 \end{bmatrix}$$

Example - element loads and Dirichlet BCs

Let us use what we have just learned and show a quick example:

- Same two-element bar model as before ⇒ stiffness matrix does not change!
- Added distributed load and non-zero prescribed displacement on the right



Element-level postprocessing

From discrete nodal displacements to continuum element fields:

Assemble and solve the global system of equations:

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From the ODE solution, recover relevant fields as function of u^e:

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Finally, plot the results! Works for displacements and any internal field (e.g. moments)

Matrix method versus FEM

The two methods give the same results for bars. However:

- The matrix method solves the strong form ODEs exactly
- FEM solves the weak form problem on the shape function space
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Then why don't we just use the matrix method for everything?

- Gluing elements through equilibrium only works in 1D
- Exact ODE solutions in 2D generally do not exist

Example – 3D frame with torsion

Full solution by hand to demonstrate all steps:

- Definition of a new element (torsion)
- Element reduction for tractability (bending)
- Element loads, support reactions, postprocessing

Values for numerical calculation:

- $EI = 1000 \text{ kNm}^2$
- $GI_t = 800 \text{ kNm}^2$
- *ℓ* = 2 m
- T = 4 kNm
- q = 6 kN/m
- m = 2 kNm/m

